

*Ollscoil na hÉireann, Gaillimh*  
*National University of Ireland, Galway*

GX 2173

**Semester II Examinations, 2003/2004**

Exam Code(s)	3BS3,3BS9
Exam(s)	Third Science, Fourth Science
Module Code(s)	MP303
Module(s)	Quantum Mechanics (Pass)
Paper No	1
Repeat Paper	
External Examiner(s)	Professor Brian Straughan
Internal Examiner(s)	Dr. Mícheál Ó Confhaola
	Dr. Michael Tuite
<b>Instructions:</b>	Full marks for <b>FIVE</b> correctly answered questions.
Duration	3hrs
No. of Answer books	
<b>Requirements</b>	
Handout	
MCQ	
Statistical Tables	Yes - Log Tables
Graph paper	
Log Graph Paper	
Other Material	
No. of Pages	4
Department(s)	Mathematical Physics

1. a. Describe the probabilistic interpretation of a wave function  $\Psi(x, t)$  for a 1dimensional particle. How is the expectation of an observable  $Q(x, \hat{p})$  defined?
- b. The probability current is defined to be  $J(x, t) = \frac{i\hbar}{2m} (-\frac{\partial \Psi^*}{\partial x} \Psi + \Psi^* \frac{\partial \Psi}{\partial x})$ . Show that
 
$$\frac{d}{dt} \text{prob}(a \leq x \leq b) = J(b, t) - J(a, t).$$
- c. Let  $\{\psi_n(x), n = 1, 2, \dots\}$  be an orthonormal set of energy eigenfunctions with eigenvalues  $E_n$ . Let  $\Psi(x, t)$  be a general wave function with initial form  $\Psi(x, 0) = \sum_{n \geq 1} c_n \psi_n(x)$ . Show that

$$\Psi(x, t) = \sum_{n \geq 1} c_n \psi_n(x) \exp(-\frac{iE_n t}{\hbar}).$$

- d. Hence show that

$$\langle \hat{H} \rangle = \sum_{n \geq 1} E_n |c_n|^2.$$

2. Consider a particle of mass  $m$  moving along a line with position  $x$  for  $0 \leq x \leq \pi$  in an infinite square well with potential

$$V(x) = \begin{cases} 0, & 0 \leq x \leq \pi \\ \infty, & \text{otherwise} \end{cases}$$

- a. Show that the normalized energy eigenfunctions are

$$\psi_n(x) = \sqrt{\frac{2}{\pi}} \sin(nx),$$

with energy  $E_n = n^2 \hbar^2 / 2m$ .

- b. Find the uncertainties  $\sigma_x, \sigma_p$  for the ground state  $\psi_1(x)$ .
- c. Show that the Heisenberg uncertainty principle  $\sigma_x \sigma_p \geq \frac{\hbar}{2}$  is satisfied.
3. Consider the 1-dimensional Harmonic Oscillator system of mass  $m$  with Hamiltonian

$$\hat{H} = \frac{\hat{p}^2}{2m} + \frac{m\omega x^2}{2}.$$

Define the ladder operators

$$a = \frac{1}{\sqrt{2m}} (i\hat{p} + m\omega x), a^\dagger = \frac{1}{\sqrt{2m}} (-i\hat{p} + m\omega x).$$

- a. Show that  $aa^\dagger = \hat{H} + \frac{\hbar\omega}{2}$  and  $a^\dagger a = \hat{H} - \frac{\hbar\omega}{2}$ .
- b. Show that with  $\xi = x\sqrt{m\omega/\hbar}$  then

$$a = \sqrt{\frac{\hbar\omega}{2}} \left( \frac{d}{d\xi} + \xi \right), \quad a^\dagger = \sqrt{\frac{\hbar\omega}{2}} \left( -\frac{d}{d\xi} + \xi \right).$$

- c. Show that if  $\psi$  is an energy eigenfunction with energy eigenvalue  $E$  then  $a^\dagger \psi$  and  $a\psi$  are eigenfunctions with energy  $E + \hbar\omega$  and  $E - \hbar\omega$  respectively. Hence show that the energy eigenvalues are  $E_n = (n + \frac{1}{2})\hbar\omega$  for  $n = 0, 1, 2, \dots$
- d. Show that the ground state energy is  $E_0 = \frac{1}{2}\hbar\omega$  with normalized eigenstate

$$\psi_0 = \left( \frac{m\omega}{\pi\hbar} \right)^{1/4} \exp(-\frac{1}{2}\xi^2).$$

- 4.
- Show that the energy eigenstates for a 1-dimensional system with even potential function  $V(x)$  can be chosen to be either even or odd functions in  $x$ .
  - Consider the energy eigenstates for a particle of mass  $m$  moving in the potential well

$$V(x) = \begin{cases} -\frac{\hbar^2}{m} < 0, |x| < \frac{\pi}{4}, \\ 0, |x| > \frac{\pi}{4}. \end{cases}$$

Show that this system has a ground state with energy  $E = -\frac{\hbar^2}{2m}$ .

- 5.
- Show for a free particle moving along a line that the energy eigenvalues are  $E_k = \hbar^2 k^2 / 2m$  for any real  $k$  for eigenfunction

$$\psi_k(x) = \exp(ikx).$$

Show that  $\psi_k(x)$  is an eigenfunction of the momentum operator  $\hat{p}$  and explain why such momentum eigenstates are not physically possible.

- A free particle (at some time  $t$ ) is described by a Gaussian wave packet

$$\psi(x) = \frac{1}{(2\pi a^2)^{1/4}} \exp\left(-\frac{x^2}{4a^2}\right) \exp(ik_0 x).$$

for some constants  $a > 0$  and  $k_0$ . Show that the uncertainty in position  $\sigma_x = a$ .

- Show that the momentum space wave function  $\phi(k) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \psi(x) \exp(-ikx) dx$  is given by

$$\phi(k) = \frac{1}{(2\pi b^2)^{1/4}} \exp\left(-\frac{(k - k_0)^2}{4b^2}\right),$$

for  $b = 2/a$ . How are the uncertainties  $\sigma_x$  and  $\sigma_p$  related in this case?

- 6.
- Prove the Schwarz inequality for states  $f, g$  in a Hilbert space

$$|\langle f, g \rangle| \leq \|f\| \cdot \|g\|,$$

where  $\|f\| = \langle f, f \rangle$  etc.

- Show for a given wavefunction  $\psi$  that  $\sigma_x^2 = \|f\|^2$  for  $f = (x - \langle x \rangle)\psi$  and  $\sigma_p^2 = \|g\|^2$  for  $g = (\hat{p} - \langle \hat{p} \rangle)\psi$
- Hence prove the Uncertainty Principle for  $x$  and  $\hat{p}$

7. Consider a particle of mass  $m$  moving in an infinite spherical potential well with potential

$$V(r) = \begin{cases} 0, & 0 \leq r < a \\ \infty, & r > a \end{cases}$$

- Show that the s-wave energy eigenfunctions are given by

$$\psi_n(r) = \frac{1}{\sqrt{2\pi a}} \frac{1}{r} \sin\left(\frac{n\pi r}{a}\right).$$

- Find the energy eigenvalues.
- Show that the probability of finding  $m$  in the region  $r < a/2$  is one half.

You may assume that the Laplacian in spherical polar coordinates is given by

$$\nabla^2 = \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2}{\partial \phi^2}.$$

8. Consider the angular momentum operators  $L_1, L_2, L_3$  obeying the commutation relations

$$[L_i, L_j] = i\hbar \epsilon_{ijk} L_k.$$

Define the ladder operators  $L_{\pm} = L_1 \pm iL_2$ .

- a. Show that the following commutation relations hold

$$[L^2, L_3] = 0,$$

$$[L_3, L_{\pm}] = \pm \hbar L_{\pm}.$$

- b. Show that  $L^2 = L_+ L_- + L_3^2 \mp \hbar L_3$ .

- c. Hence show that the  $L^2$  and  $L_3$  eigenstates  $|lm\rangle$  obey

$$L^2 |lm\rangle = \hbar^2 l(l+1) |lm\rangle,$$

$$L_3 |lm\rangle = \hbar m |lm\rangle,$$

for  $l = 0, \frac{1}{2}, 1, \frac{3}{2}, \dots$  with  $m = -l, -l+1, \dots, l-1, l$ .

- d. Comment on the physical realization of such eigenstates for integer and half-integer  $l$ .

9. A two particle system consists of spin  $s_1 = \frac{3}{2}$  particle together with a spin  $s_2 = \frac{1}{2}$  particle. Describe the total spin eigenstates  $|sm\rangle$  of the system in terms of the basis  $\{|s_1 m_1\rangle |s_2 m_2\rangle\}$  and hence find the Clebsch-Gordan coefficients  $C_{m_1 m_2 m}^{s_1 s_2 s}$ .

You may assume the relation  $L_{\pm} |lm\rangle = \hbar \sqrt{l(l \mp m)(l \pm m + 1)} |lm \pm 1\rangle$  for orthonormal simultaneous eigenstates of  $L^2$  and  $L_3$ .

#### Useful formulae

$$\hat{p} = -i\hbar \frac{\partial}{\partial x}, \text{ 1-d Momentum Operator}$$

$$\hat{H} = \frac{1}{2m} \hat{p}^2 + V(x), \text{ 1-d Hamiltonian Operator}$$

$$i\hbar \frac{\partial}{\partial t} \Psi(x, t) = \hat{H} \Psi(x, t)$$

$$= -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \Psi(x, t) + V(x) \Psi(x, t), \text{ 1-d Schrödinger Equation}$$